

# AGN Jets

A deep-field astronomical image showing a bright, central point source (the AGN) with a prominent, elongated, and slightly curved jet extending upwards and to the left. The jet is composed of numerous smaller, reddish-pink spots, suggesting a complex structure of gas and dust. The background is a dense field of smaller, fainter stars or galaxies, also appearing as reddish-pink spots.

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Credit: NASA/CXC/A.Siemiginowska(CfA)/J.Bechtold(U.Arizona)

# Outline

- ◆ Before Chandra
- ◆ Physics of Jets · What we expect
  - ◆ Emission Processes
  - ◆ Types of Jets
  - ◆ Surrounding medium
- ◆ Detection · How we see them
- ◆ Results · What we see
- ◆ Future- what we'd like to see

# Before Chandra

- ◆ Jets mainly studied in radio · Xray jets thought „exotic“
- ◆ Einstein and ROSAT had few · M87, Cen A, NGC 6251
- ◆ Chandra first light · Xray jet PKS 0637-752
- ◆ Xrays found where no optical emission
- ◆ Synchrotron model breaks down from radio to Xray
- ◆ Now it seems that ALL radio jets have Xray counterparts

# Synchrotron Review

- ♦ See Rybicki and Lightman Ch.6

$$\text{Larmor Formula: } P = \frac{(2 q^2 \dot{u}^2)}{(3 c^3)} = \frac{(2 q^2 \vec{a} \cdot \vec{a})}{(3 c^3)}$$

*For relativity:  $\vec{a} \cdot \vec{a}$  is invariant*

$$\text{Beaming: } a_{\parallel}' = \gamma^3 a_{\parallel}, \quad a_{\perp}' = \gamma^2 a_{\perp}.$$

$$a_{\perp}' = \omega_B v_{\perp}.$$

$$\omega_B = \frac{(qB)}{(\gamma mc)}$$

$$P = \frac{(2 q^2)}{(3 c^3)} \gamma^4 \frac{(q^2 B^2)}{(\gamma mc^2)} v_{\perp}^2.$$

# More Synchrotron Review

*Classical Electron radius* :  $r_0 = \frac{q^2}{(mc^2)}$

$$U_B = \frac{B^2}{(8\pi)} \quad \sigma_T = 8\pi r_0/3$$

$$\tau_{rad} = \frac{E}{(-dE/dt)}$$

*Voila* :  $-dE/dt = 2\sigma_T c \gamma^2 \beta^2 U_B \sin^2 \alpha (1)$

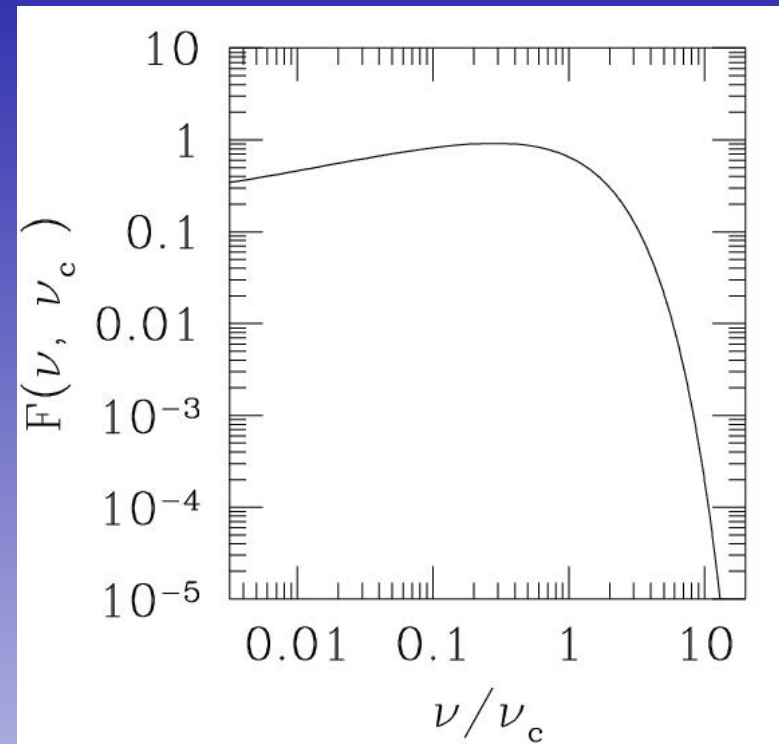
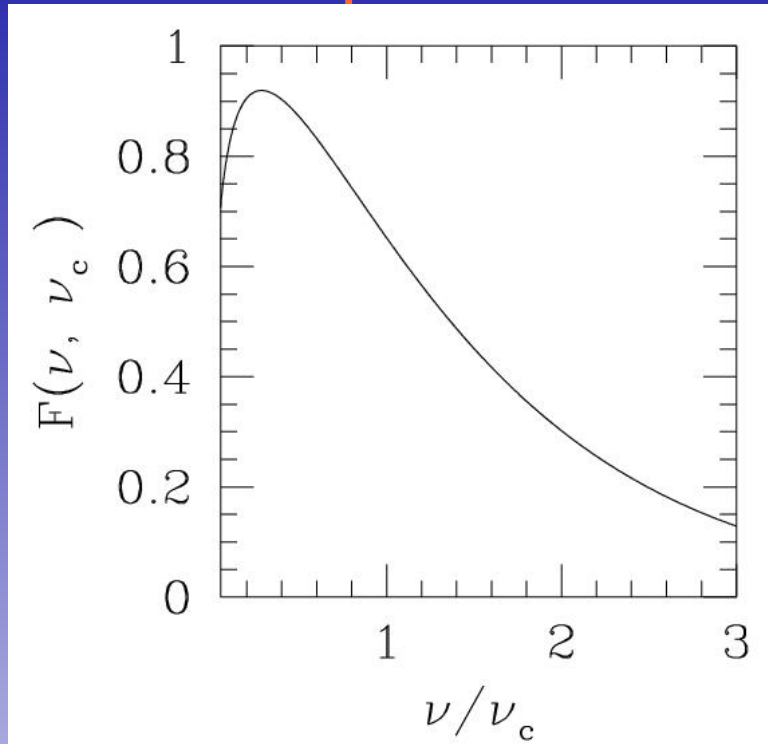


$$\nu_g = \gamma \omega_B / 2\pi$$

*Critical Freq.* :  $\nu_c = (3/2) \gamma^2 \nu_g \sin \alpha$

# Synchrotron · SDF and Luminosity

- ◆ Critical Freq is 30GHZ



From Worrall and Birkinshaw

$$X = \frac{\nu}{\nu_c}$$

$$F(\nu, \nu_c) = X \int_X^\infty (K_{(5/3)}(\xi) d\xi)$$

$$L_\nu \propto \int F(\nu, \nu_c) N(\gamma) d\gamma$$

# Synchrotron · Power Laws and Luminosity

- ◆ The Big Punchline

$$L_\nu \propto \int F(\nu, \nu_c) N(\gamma) d\gamma, \quad N(\gamma) d\gamma = \kappa \gamma^{(-p)} d\gamma$$

$$\text{Requires: } \gamma_{min}^2 \nu_g \ll \nu \ll \gamma_{max}^2 \nu_g$$

Analytical Result is a mess!!

- ◆ However,  $L_\nu \propto \nu^{(-\alpha)}$

- ◆ Physics · Pret, mean probability per cycle that particle remains in jet and is shocked again

$$\text{Ultarel. case: } \alpha = 1 + \frac{(\ln 1/P_{ret})}{(\ln E_f/E_i)} \quad (\text{Achterberg, et al 2001})$$

Fermi acceleration!

$$\text{Nonrel. case: } \alpha = \frac{(\rho_d/\rho_u + 2)}{(\rho_d/\rho_u - 1)}$$

# Phew! Onto Inverse Compton

- ♦ Exact same energy rate loss formula! Provided:
- ♦  $\gamma h \nu_o \ll m_e c^2$  – low energy photons!
- ♦ Spectral distribution function differs:

$$X = \frac{\nu}{(4\gamma^2 \nu_o)}$$

$$F(\nu, \nu_o, \gamma) = X f(x), \quad f(x) = X(1 + X - 2X^2 + 2X \ln X)$$

$$\bar{X} = \frac{(\int X f(X) dX)}{(\int f(X) dX)} = 1/3, \quad \bar{\nu} = (4/3)\gamma^2 \nu_o$$

- ♦ When Compton scattering is on synchrotron photons from jet, is synchrotron self-compton



# Thermal Bremsstrahlung Emission

Stated: 
$$\frac{dW}{(d\omega dV dt)} = \frac{(16\pi e^6)}{(3\sqrt{3}c^3 m^2 \nu)} n_e n_i Z^2 g_{ff}(\nu, w)$$

Maxwell Velocity Dist' n (Isotropic): 
$$dP \propto v^2 \exp\left(\frac{(-mv^2)}{(2kT)}\right) dv$$

For a photon of freq.  $\nu$ ,  $h\nu \leq \frac{1}{2}mv^2$

- ◆ Average  $dW(\nu, \omega)/(d\omega dV dt)$  over the velocity distribution

- ◆ 
$$\epsilon_\nu = 32 \frac{\pi}{3} \left(2 \frac{\pi}{3}\right)^{.5} Z_i^2 g_{ff} n_i n_e m_e c^2 r_o^3 \left(\frac{(m_e c^2)}{(kT)}\right)^{.5} e^{\left(\frac{(-h\nu)}{(kT)}\right)}$$



Integrate over freq., convert to P

$$\tau_{cool} = 2.5 \frac{(kT)}{(\epsilon \mu m_H)}$$

# FR classification

- ◆ Paper: Fanaroff and Riley, 1974 MNRAS · Cambridge One Mile telescope
- ◆ "There is a definite relationship between relative positions of the high and low brightness regions of radio sources and their luminosity"
- ◆ Classified by ratio: distance between highest central bright. regions on either side of nucleus and total extent · relative position of hotspots
  - ◆  $> .5$  Class II,  $< .5$  Class I, High L Class II, Low L Class I
  - ◆ Class I more complex structure

# FR I



- ♦ FR I galaxy 3C 31 · note that jets starts right near core

# FR I physics

- ◆ Low power, and a well defined beam head · this means good contact with external medium
- ◆ Based on 3C 31, jet has 3 regions: narrow jet w/ relativistic flow, flaring region bright in radio, outer region with steady deceleration
- ◆ Outer region is where deflection takes place
  - ◆ Knots are a mystery
  - ◆ Acceleration takes place in knots

# FR II



- ◆ FR II galaxy 3C 220.1 · clear separation of lobes from core · high power

# FR II Physics

- ◆ FR II is supersonic w.r.t ISM- bow shock forms at end of jet · we see „well collimated jets feeding edge brightened lobes“

- ◆ 
$$c_s = \sqrt{\left(\frac{(\gamma k T)}{(\mu m_H)}\right)}$$

- ◆ 
$$M = \frac{v_{adv}}{c_s}$$
 Assume strong shock (Rankine-Hugoniot)

Get simple relations

$$\frac{P_2}{P_1} = (5 M^2 - 1) / 4, \quad \frac{\rho_2}{\rho_1} = \frac{(4 M^2)}{(M^2 + 3)}$$

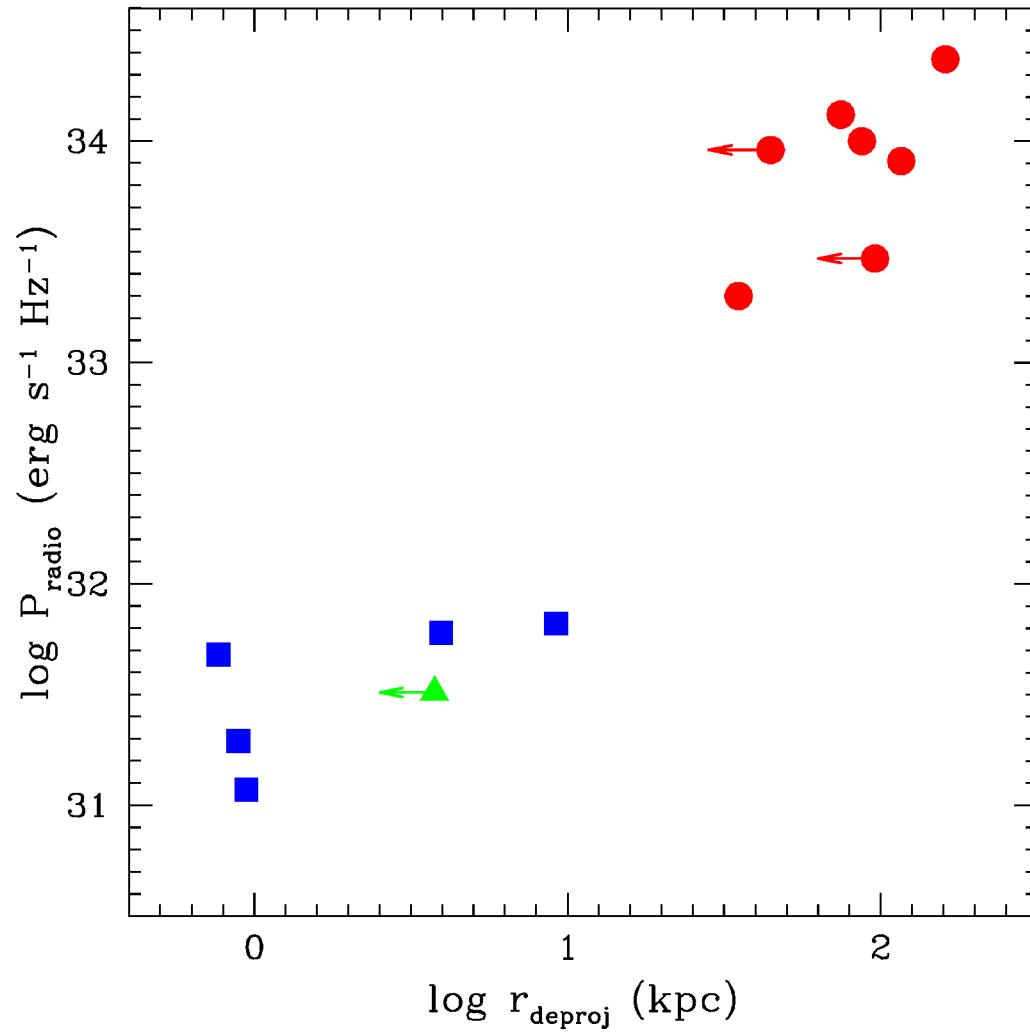
$$\frac{T_2}{T_1} = (5 M^2 - 1) \frac{(M^2 + 3)}{(16 M^2)}$$

# FR I/ II Physics

- ◆ Difference (according to Sambruna) lies in jet pressure · fits with idea of luminosity discrimination
- ◆ Jet pressure in FR II becomes comparable to ambient farther out · this is where shocks form
- ◆ Would expect innermost knot closer to core in FR I than in FR II · explore?
- ◆ In FR I knots fit radio-optical synchrotron spectrum
- ◆ FR I Xray morph. Fairly uniform

# FR I/ II Physics

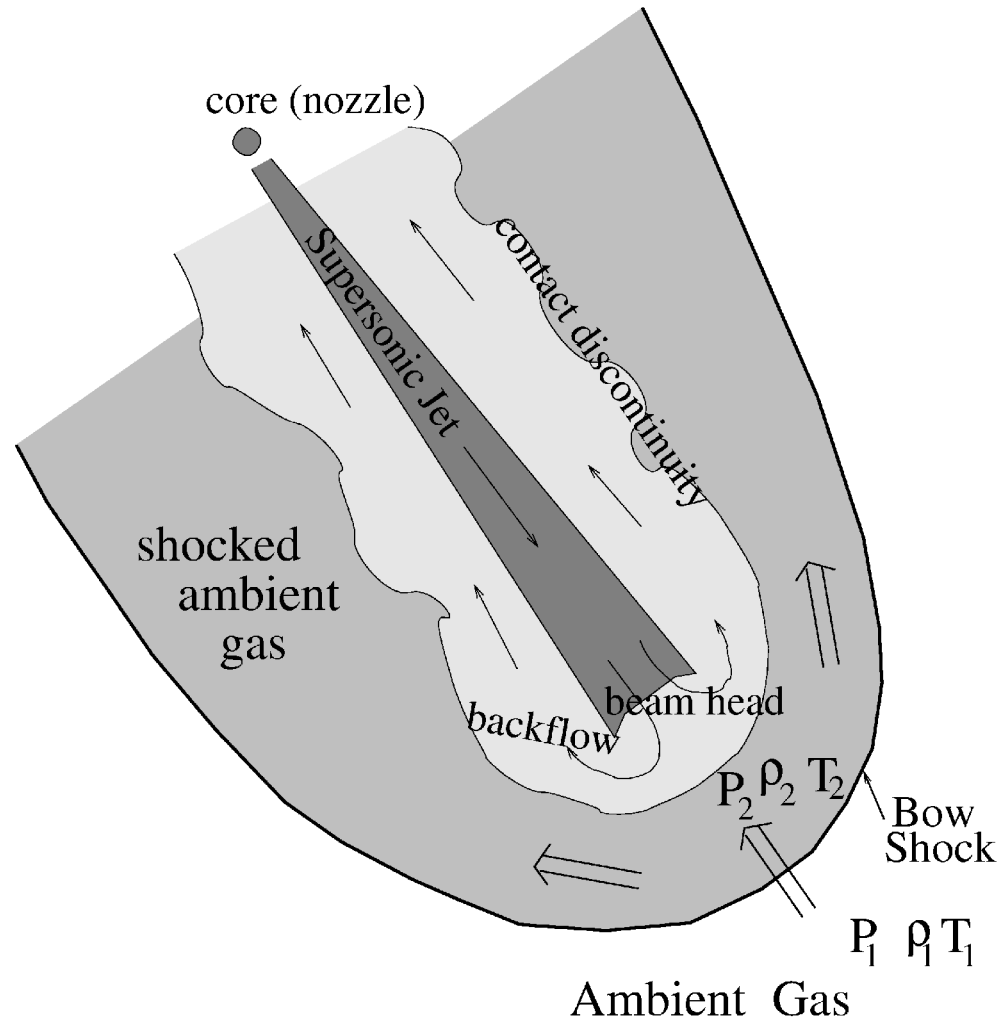
From Sambruna, et al.





# Jet Schematic

From Worrall and Birkinshaw



- ♦ Jet terminates at hotspot - light gray is radio emitting fluid · heats ambient gas producing X-rays

# Jet Structures

- ◆ Jet features defined based on radio
- ◆ Jet · narrow feature, 4x as long as it is wide
- ◆ Lobe · radio emission not in jet
- ◆ Hotspot · compact feature in lobe
- ◆ Knot · compact feature in jet
- ◆ End of jet: end of emission, 90 degree change of direction, decollimation by 2x
- ◆ Knots are most likely to have X-ray counterparts

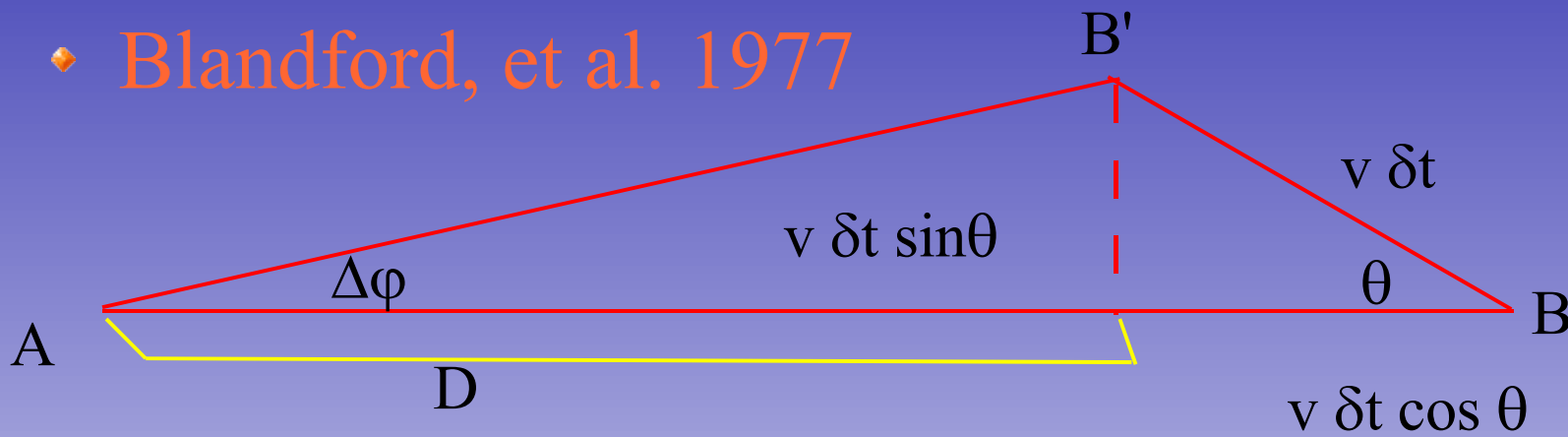
# Magnetic Fields

- ◆ Field orient. determined by linear polarization in radio
- ◆  $\pi_L = \frac{(\alpha + 1)}{(\alpha + 5/3)}$
- ◆ Significant pol. indicates lack of thermal material · otherwise differential between front and back of jet would depolarize it
- ◆ Circular polarization (mentioned later) also seen in pc scale jets
- ◆ 

Is this e+ e-?

# Superluminal Motion

- VLBI obs. of AGN show jets w/ multiple comp. around bright cores
- These comp. Expanding · separating w/  $Vt \sim 10c$
- Blandford, et al. 1977



Adapted from Peterson (1997)

$$\beta_T = \frac{v_T}{c} = \frac{D (\Delta \phi)}{c (\Delta t)} = \frac{(v \sin \theta)}{(c (1 - \beta \cos \theta))} = \frac{(\beta \sin \theta)}{(1 - \beta \cos \theta)}$$

Maximize  $\beta$ , get  $\beta_T^{max} = \beta \gamma$ ,  $\gamma$  can be arb. large

# Bulk Relativistic Motion · Effects

- ◆ Relativistic Doppler Factor:

- ◆ 
$$\delta = \frac{1}{(\Gamma(1 - \beta \cos \theta))}$$
- ◆ 
$$R_J = \left( \frac{(1 - \beta \cos \theta)^{-(\alpha+2)}}{(1 + \beta \cos \theta)} \right)$$
- ◆  $L_{nu}$  increases by  $\delta^{(2+\alpha)}$
- ◆ Boosting by factor of alpha due to blueshift

- ◆ Core dominance:

$$R_{cd} = \frac{R_{(cd, (\cos \theta = 0))}}{2} \left( (1 - \beta \cos \theta)^{-(\alpha+2)} + (1 + \beta \cos \theta)^{-(\alpha-2)} \right)$$

# External Medium

- ♦ Jets mostly distinct from nucleus · X-rays arise from interaction of hot jet gas with cool dense medium

- ♦ Jet is in hydrostatic equilibrium · like stellar interior

- ♦ Ideal Gas:  $\nabla P = -\rho \nabla \phi$   
 $P = \frac{(n_p k T)}{(X \mu)}$ , solar abund.  $\mu = .6$

- ♦ Spherical sym.:  $\frac{(d\Phi)}{(dr)} = \frac{(G M_{tot}(r))}{r^2}$

# More External Medium

- ◆ Follows "Isothermal Beta Model" - see Cavaliere and Fusco-Fermiano, 1978

- ◆  $\rho = \rho_0 \left(1 + \frac{r^2}{r_c^2}\right)^{-1.5\beta}$
- ◆  $\beta = 2/3 \frac{(\mu m_h)}{(kT)} \frac{(GM_c)}{r_c}$

- ◆ Beta depends on ratio of grav. pot. to thermal E per unit mass
- ◆ NFW say its unphysical · measured beta's don't match theory · NFW ain't perfect either
- ◆ Bottom line · density has power-law structure with dimensionless parameter representing balance of gravity and thermal motions

# More External Medium -Surface Brightness and Ion Ratio

- ◆ Density and T of medium can be determined from surface brightness, but it's a mess
- ◆ XSPEC uses the magic normalization mentioned in class

- ◆ 
$$N = \frac{\left( \int n_e n_p dV \right)}{\left( 4 \pi D_A^2 (1+z)^2 \right)} = \frac{\left( (1+z) S_\nu \right)}{\Lambda_\nu}$$

- ◆  $D_A$  is angular diameter distance and  $(1+z)$  is for expanding universe,  $\Lambda$  is emissivity

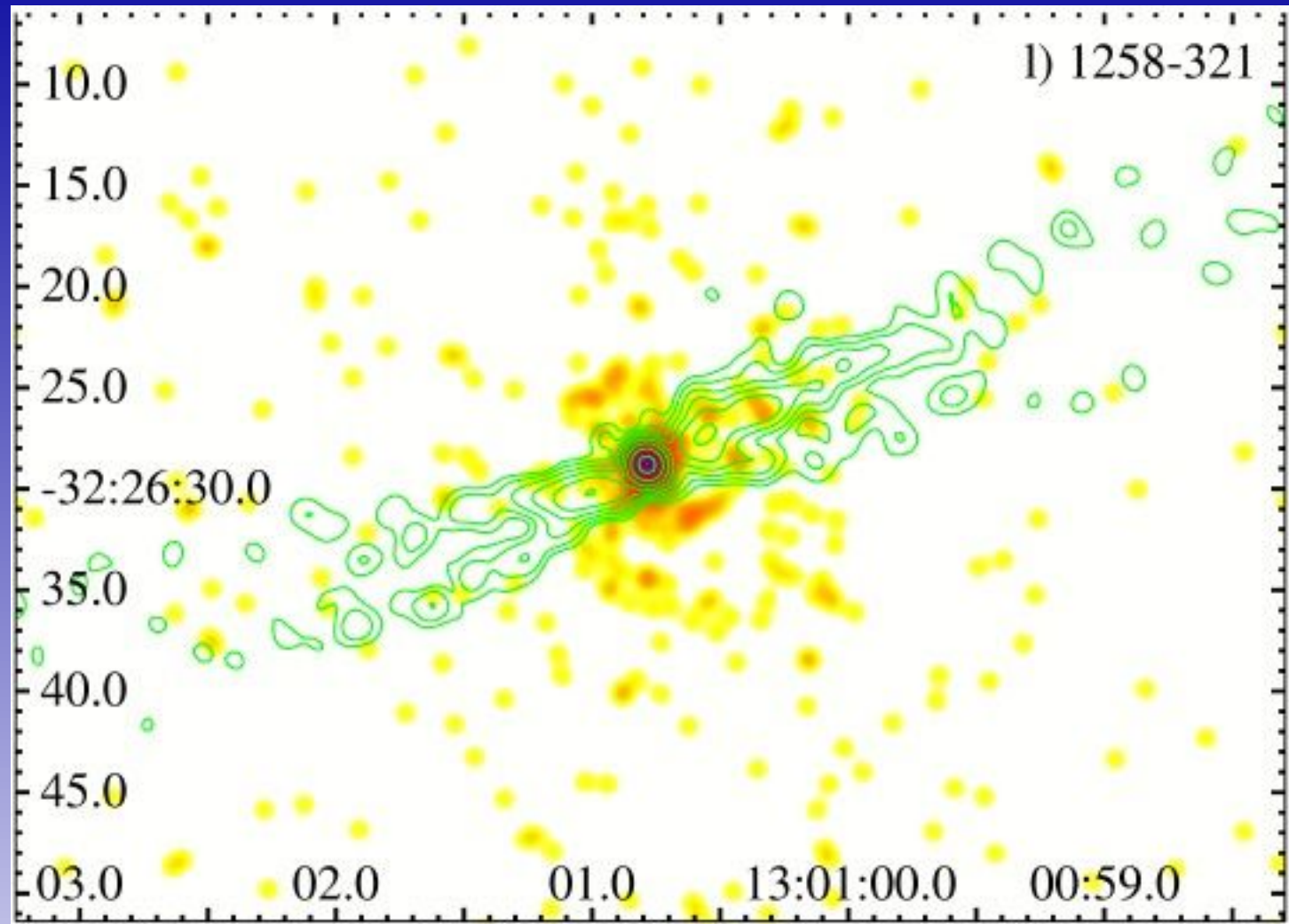
- ◆ S and  $\Lambda$  are fit,  $n_e$  and  $n_p$  are "pulled out" using isothermal beta, and

$$\frac{n_e}{n_p} = \frac{(1+X)}{(2X)} = 1.8 \text{ for solar}$$



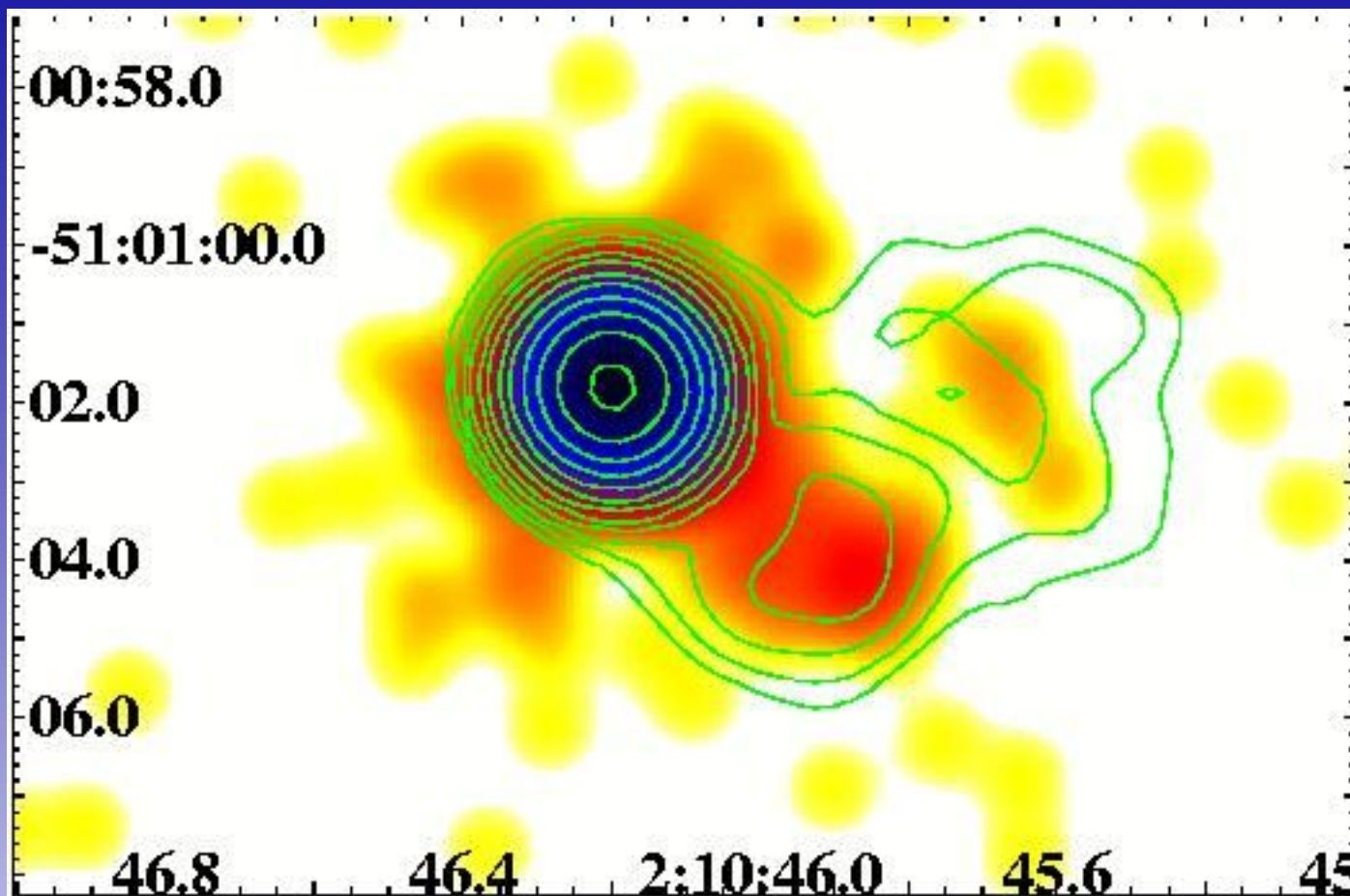
# PKS 1258-321

Marshall et al.



- ◆ Core dominated, FRI, no X-ray emission from the jets
- ◆ Extended X-ray halo?

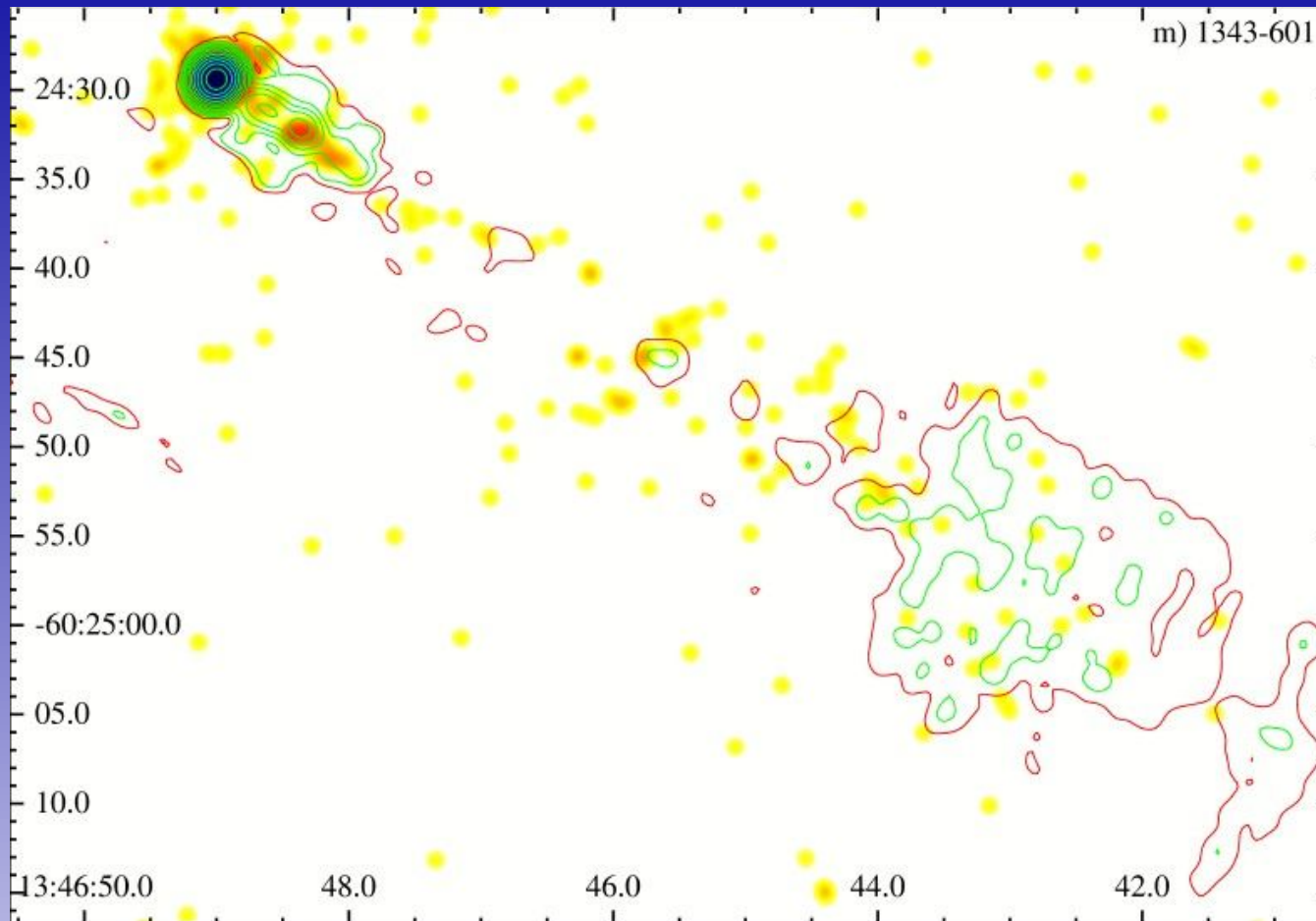
# PKS 0208-512



Marshall et al.

- ◆ 4% long jet in X-ray, coincident with radio · X-ray cuts out at 90 degree bend · radio polarization is at 90 degrees to jet direction

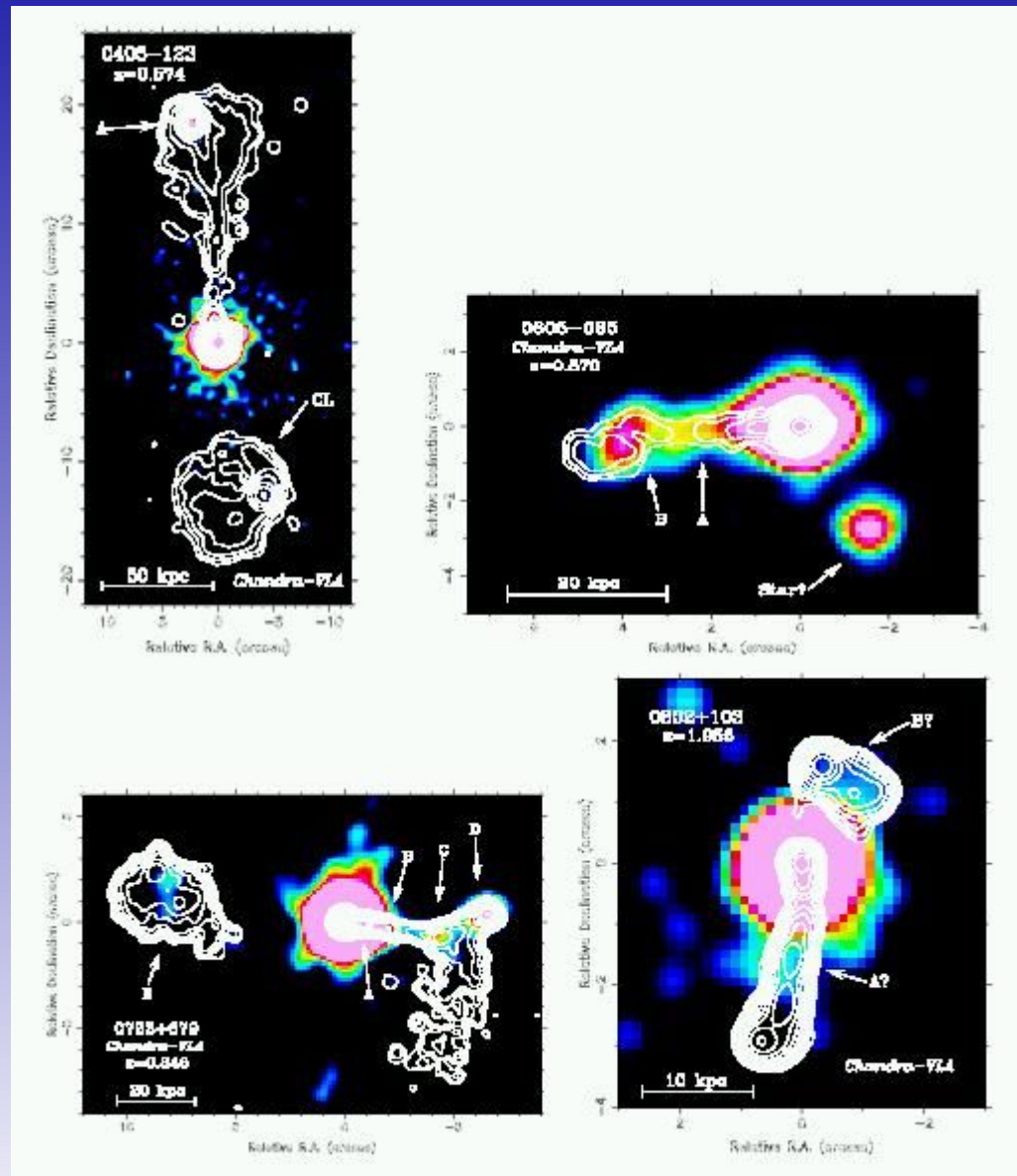
# PKS 1343-601 / CEN-B



Marshall et al.

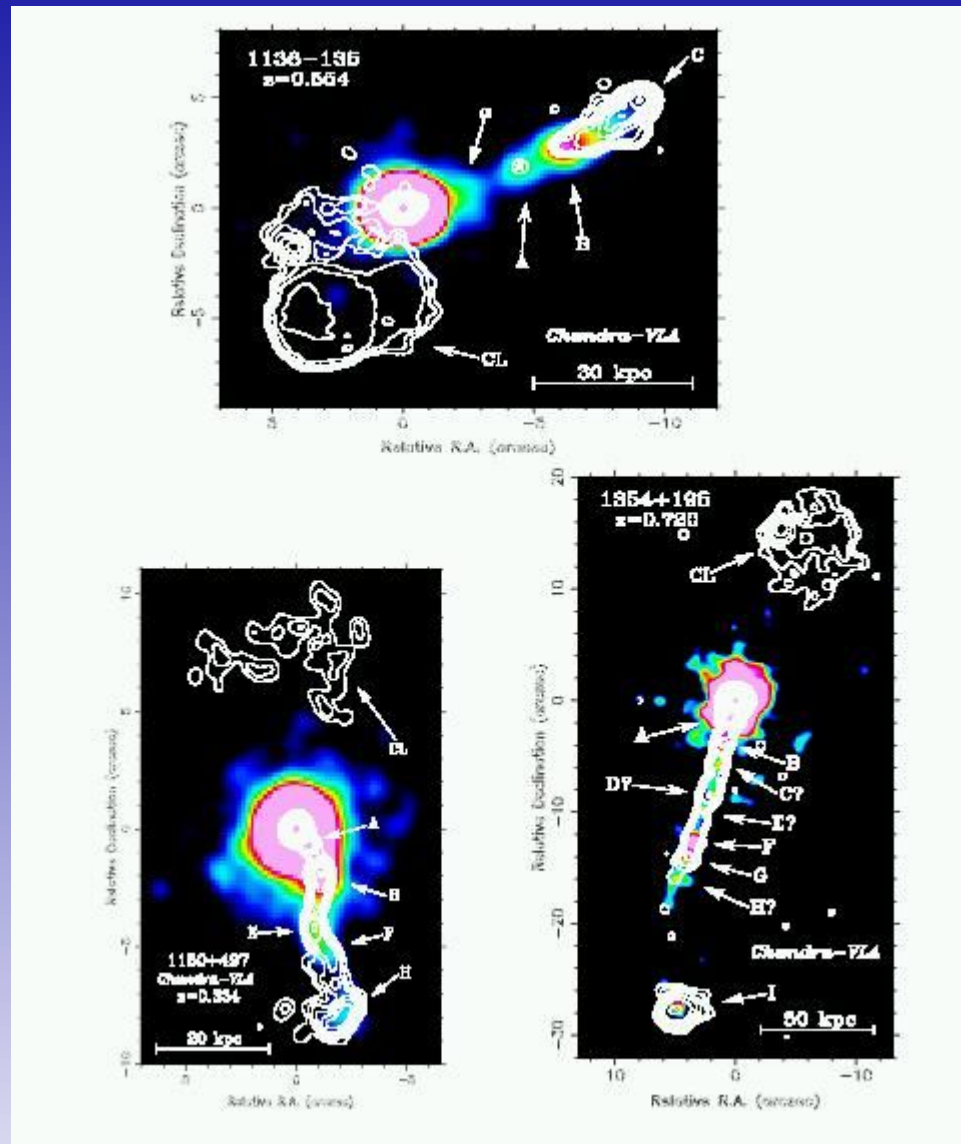
- ♦ Very one sided, knot detected in X-ray, some yellow blobs significant, upstream and in lobe

# Sambruna et al.





# Sambruna et al.



# What can we learn?

- ◆ Sambruna et al find  $\sim 59\%$  detection in X-ray of radio features, around same in optical
- ◆ Seems that most all FR I's have X-ray counterparts
- ◆ All sources in Sambruna sample are one-sided - beaming
- ◆ Composition:  $e^-$  &  $e^+$  pairs expected near AGN ·  
Circular polarization from mode conversion detected
- ◆ Evidence for  $p^+$  &  $e^-$  plasma exists too ·  $p^+$  outnumber  $e^-$  in CR by 100x
  - ◆ Also, interaction w/ medium should enhance  $p^+$  content
  - ◆ However, data point to a low mass jet, insufficient pressure for a  $p^+$  &  $e^-$  jet

# Spectral Modeling

- ◆ SED modeled for aperture of 1-1.5", limited by Xray
- ◆ SED power law from radio through optical to Xray determines emission mechanism (IC, synch., thermal)
- ◆ If optical lies on power law between radio and Xray, then 1 method works for all
- ◆ If below, then synch. provides low freq. and IC provides high
- ◆ Synchtron dominates when:  $\alpha_{ro} < \alpha_{ox}$
- ◆ For most jets, Xray is a separate process
- ◆ Alpha rx increases from inside of jet out
- ◆ Confirms theory of W&B

# More Spectral Modeling

- ◆ A single e- population can produce both synch. and IC present · need input photons from CMB
- ◆ Any seed photons will do · CMB is most likely
- ◆ Energy density of CMB scales  $(1+z)^4$
- ◆ For knots inside host, starlight works too
- ◆ As shown earlier, relative lifetimes for the emission processes vary, synch. short @ short WL
- ◆ Restricts synch. to more compact reg.
- ◆ IC/CMB model · huge cooling times
- ◆ Need consistent Xrays along jet - PKS 1954+196 shown before



# More Spectral Modeling

- ◆ Sambruna finds trends based on large sample:
- ◆ "Common conditions along jet determine the observed properties of emission features"
- ◆ Possible reason for FR I/II: location where jet and ambient P equalize · causes major slowdown
- ◆ Discrim. Param · Jet P (power), more luminous sources slow down later
- ◆ "Xray emission from extended radio jets with  $\alpha_{rx}$  around .8 is common in AGN"
- ◆ Knots have different phys. process

# More Spectral Modeling

- ◆ Marshall has similar findings on IC/CMB, interesting results on PA
- ◆ First, need to define min. B field strength · find energy in e- by integrating power law (at min.  $U_e = 4/3 U_b$ ):

$$U_e = \kappa m_e c^2 \int_{(\gamma_{min})}^{(\gamma_{max})} \gamma \gamma^{(-p)} d\gamma = \kappa m_e c^2 \int (\gamma_{min})^{(\gamma_{max})} \gamma^{(-2\alpha)} d\gamma$$

- ◆ Know  $L_\nu \propto \kappa \nu^{-\alpha} B^{(\alpha+1)}$
- ◆ Integrate and plug in for kappa, define K as baryon ratio, eta=f=filling factor

$$U_{part.} = C_1 B^{-(\alpha+1)} \frac{(1+K)}{(\eta V)} L_{nu} f(\alpha)$$

◆ Marshall: alpha=.8, gamma=1

# More Spectral Modeling

- ◆ From there, derived quantities and some assumptions lead to:

- ◆  $\Gamma \delta (1 + u_j') = \frac{(1 - \beta + \mu - \mu \beta)}{(1 - \mu \beta)^2}$ ,  $\mu_j = \text{cosine of } l - o - s \text{ angle}$

- ◆ Assume  $\Gamma = 10$ , know  $\delta = [\Gamma (1 - \mu \beta)]^{-1}$

- ◆ Can solve for PA!

- ◆ Jet angles around 20 to 30 degrees

# What We Don't know:

- ◆ Lots! Biased samples · current surveys have selection effects:  $L$ ,  $z$
- ◆ Need obs. of Xray ISM
- ◆ Deep observations · would like to see continuous emission along all jets
- ◆ Limited to 1% aperture · could be overestimating sizes
- ◆

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