

Neutrino astrophysics

The Fermi acceleration of electrons producing the observed **GRB** or AGN γ -ray, X-ray and optical photons is likely also to accelerate protons, either in internal shocks [1] or external shocks [4]. The extreme electromagnetic brightness implies a high intra-source density of photons, and a high optical depth to $p\gamma$ photo-meson interactions. Interactions between protons of $E_p \gtrsim 10^{16}$ eV and photons of ϵ_γ leads to neutrinos whose energy in the observer frame is $E_\nu, \text{ GeV} \sim 10^{-2}\Gamma^2/\epsilon_\gamma, \text{ GeV} \sim 100$ TeV and extending to PeV for a GRB emitting MeV photons from a jet of Lorentz factor $\Gamma \sim 10^2$. A comprehensive discussion of Fermi acceleration and neutrinos in GRBs and AGNs was given in [3].

A different mechanism for neutrino production in GRB [6] is inelastic p, n collisions in the outflow, caused by the relative drift between entrained neutrons and protons when they decouple; these neutrinos have energies $E_\nu \sim 40 - 50$ GeV. Another neutrino component [7] is expected from core-collapse “long” GRBs, when internal shocks occur while the jet is still inside the star. Photo-meson and pp neutrinos at $E_\nu \sim 1$ TeV arise then as precursors $\sim 10 - 100$ s before the γ -rays from the post-emergence jet. There may be a large number of ‘failed’ GRB, where **choked jets** did not emerge, leading to an orphan neutrino burst [7] without a corresponding electromagnetic signal. Neutrinos would serve as ‘tomographical’ probes of the stellar structure of the stellar progenitors [8]. More detailed predictions for choked and sub-photospheric jet neutrinos were given in [9], and a different, proton-neutron converter mechanism for neutrino production in buried radiation-dominated flows was discussed in [10].

The shocks in the jets of **AGNs** may also be photo-meson neutrino sources. However, there are many more AGNs with modest or no jets, and the cores of such ‘radio-quiet’ AGNs may be promising neutrino sources. This is due to shocks between blobs in the inner accretion flow or the incipient jets near the central black hole [11, 12]. **Hypernovae**, or jet-powered supernovae, are core-collapse supernovae exhibiting semi-relativistic ejecta velocities, and these can accelerate protons or ions to relativistic energies leading to GeV to PeV neutrinos [13, 14]. **Magnetars** are ultra-high magnetic field ($\gtrsim 10^{14}$ Gauss) neutron stars, which give rise to soft-gamma repeaters (SGRs), burst-like events due to magnetic reconnection leading to a fireball where shocks Fermi accelerate electrons as well as protons, which can lead to photo-meson neutrinos [15] or curvature radiation neutrinos [16]. In newly formed magnetars, wake-field acceleration can accelerate protons which interact with the SN ejecta leading to significant neutrino fluxes [17].

The diffuse TeV-PeV neutrino background observed by *IceCube*, e.g. [18], appears to be due to so-far unidentified sources at cosmological distances. Since $p\gamma$ and pp interactions produce both VHE neutrinos and γ -rays comparable numbers, the diffuse γ -ray background observed by the *Fermi* satellite imposes a strong constraint, since the latter is largely due to blazars, which however do not correlate with the neutrinos. Most optically thin sources would overproduce γ -rays, but γ -ray-dark sources such as low-luminosity or choked GRBs [19], or shocks from high-redshift merging galaxies and clusters [20] are likely sources which satisfy the constraints. A wider discussion is given in a recent review [21].

References

- [1] Waxman, E, 1995, PRL, 75:386
- [2] Waxman, E & Bahcall, JN, 1997, PRL, 78:2292
- [3] Rachen, J & Mészáros, P, 1998, PRD, 58:123005
- [4] Böttcher, M, Dermer, C, 1998, ApJ, 499:L131
- [5] Waxman, E & Bahcall, JN, 1999, PRD, 59:023002
- [6] Bahcall, JN, Mészáros, P, 2000, PRL, 85:1362
- [7] Mészáros, P and Waxman, E, 2001, PRL 87:171102
- [8] Razzaque, S, Mészáros, P and Waxman, E, 2003, PRD 68:3001
- [9] Murase, K.; Kashiyama, K.; Meszaros, P., 2013, Phys.Rev.Lett., 111:131102
- [10] Kashiyama, K., Murase, K. and Meszaros, P., 2013, Phys.Rev.Lett., 111:131103
- [11] Alvarez-Muñiz, J and Mészáros, P, 2004, Phys.Rev.D, 70:123001
- [12] Pe’er, A., Murase, K., and Meszaros, P., 2009, Phys.Rev. D, 80:123018
- [13] Wang, X-Y, Razzaque, S, Meszaros, P and Dai, Z-G, 2007, Phys.Rev.D, 76:083009
- [14] Senno, N., Meszaros, P., Murase, K., Baerwald, P. and Rees, M.J., 2015, ApJ, 806:24
- [15] Ioka, K, Razzaque, S, Kobayashi, S, Mészáros, P, 2005, ApJ, 633:1013
- [16] Herpay, T., Razzaque, R., Patkos, A. & Mészáros, P., 2008, JCAP, 8:025
- [17] Murase, K., Mészáros, P. and Zhang, B., 2009, PRD, 79:103001
- [18] Gaisser, T. and Halzen, F., 2014, Annu.Rev.Nuc.Part.Sci., 64:101-123
- [19] Senno, N., Murase, K. and Meszaros, P., 2016, Phys.Rev.D, 93, 083003
- [20] Yuan, C.C., Meszaros, P., Murase, K. and Jeong, D., 2018, ApJ, in press (arXiv:1712.09754)
- [21] Meszaros, P., 2017, Annu.Rev.Nuc.Part.Sci., 67:45-67